

Study of the lithium–nitrogen–hydrogen system

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Abstract

We investigated lithium–nitrogen–hydrogen compounds as potential hydrogen storage materials and report consistent results for the amount of hydrogen gas sorbed, weight change, and phase content for our various samples. We observed a sluggish and incomplete reaction between hydrogen and lithium nitride of >7 wt.% hydrogen uptake at 160 °C with no observable hydrogen release into vacuum, and a ~5.2 wt.% cyclable hydrogen sorption at 240 °C consistent with a reversible lithium amide/lithium imide reaction. The plateau pressure for the amide/imide reaction is quite low, ~10 kPa at 240 °C.

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Using hydrogen-powered fuel cells for automotive transportation is an exciting conceptual solution for two major societal concerns: fossil fuel dependence and pollution. Implementation, however, requires technological breakthroughs, particularly regarding efficient production of large quantities of hydrogen from non-fossil fuel sources and practical on-board hydrogen storage. Considerable research is focused on new solid substances for hydrogen storage that include metal hydrides, physisorptive materials, and hydrogen compounds. So far it has not been possible to achieve all desirable material properties for hydrogen storage systems simultaneously in a single substance [1]. We report our research on new potential hydrogen storage materials in the lithium–nitrogen–hydrogen system [2–7].

We obtained lithium nitride (Li₃N) powder from Alfa Aesar, and our chemical analysis shows that it is 51 wt.% lithium (pure Li₃N is 59.8 wt.% lithium). Our X-ray diffraction (XRD) analysis reveals it to be a mixture of the room temperature form of Li₃N and the metastable high-pressure form (As₃Na-type), and that the major impurity phases are

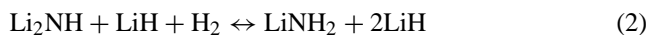
LiOH and Li₂O. We perform hydrogen (H₂) sorption experiments on as-received Li₃N powder using XRD, volumetric, and gravimetric techniques. For real-time in situ XRD experiments, we use Cu K α radiation in a Bruker AXS General Area Detector Diffractometer System. We load and seal samples into 1 mm quartz capillary tubes and mount them in a sample cell and furnace, described previously [8], in an Ar gas glove box. We collect XRD data using an 0.8 mm collimator, 15 cm detector distance, goniometer angles of $\Omega = 20^\circ$ and $2\theta = 34^\circ$, and 1–5 min integration times. This allows collection of XRD data for angles $16^\circ < 2\theta < 51^\circ$. For isothermal volumetric experiments, we use a commercial Gas Reaction Controller, described previously [9,10], and measure equilibrium pressures and temperatures to accurately determine amounts of gas in the various chambers of the apparatus. For gravimetric experiments, we use a Cahn Model 2151 high-pressure thermogravimetric analyzer (TGA) [11,12]. Carefully weighed ~300 mg powder samples are placed into quartz sample buckets and covered with pentane to protect them from exposure to air when removed from the glove box for loading into the TGA. The TGA is thoroughly purged with He gas while the pentane evaporates completely. Hydrogen sorption is directly measured as weight changes after correct-

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ing for non-gravitational forces on the electrobalance due to buoyancy, gas flow velocity, and electrobalance temperature.

We reacted Li_3N with H_2 according to the equations:



These reactions are sufficiently exothermic that, with the excess H_2 typically present in our high-pressure experiments, they proceed to the right and ultimately toward the completion of Eq. (2). The resulting mixture of LiNH_2 and LiH in a 1:2 molar ratio is consistent with our in situ XRD, volumetric, and TGA experiments presented below. Chen et al. [2] have shown that full reversal to Li_3N requires high temperatures, hence the conditions of the reaction described by Eq. (2) are more compatible with practical reversible H_2 storage.

We performed in situ XRD measurements on Li_3N initially in 550 kPa H_2 and then heated to $\sim 160^\circ\text{C}$. The temperature was then increased in 15°C increments over 1800 min to 235°C . At ~ 1860 min, the H_2 was evacuated using a mechanical vacuum pump for ~ 960 min. Finally, at ~ 2850 min, the sample holder was pressurized to 625 kPa H_2 . The in situ XRD data for Li_3N contain more than 1000 diffraction patterns collected during this ~ 50 h experiment, and Fig. 1 shows representative patterns resulting from hydriding and de-hydriding Li_3N . The room temperature diffraction pattern shows a weak but detectable diffraction peak at $2\theta \approx 30.5^\circ$ that is not observed prior to the initial H_2 pressurization, and we attribute it to a reaction of H_2 with Li_3N forming LiNH_2 and/or lithium imide (Li_2NH). During hydriding, the diffraction pattern at 2 min and $\sim 160^\circ\text{C}$ shows increasing peak intensity for LiNH_2 , and at 1800 min and 235°C shows increasing peak intensity for LiNH_2 and LiH and decreasing peak intensity for Li_3N , consistent with Eqs. (1) and (2). During de-hydriding, the diffraction pattern at 2400 min and 253°C shows a mixture of Li_2NH and LiH . Because LiNH_2 and Li_2NH have very similar diffraction patterns due to very similar crystal structures, the strongest diffraction peak for each phase has a nearly identical position at $2\theta \approx 30.5^\circ$, and consequently the intensity of this peak remains relatively constant during de-hydriding. We monitor the decomposition of LiNH_2 into Li_2NH according to the reverse of Eq. (2) by observing the disappearance of low intensity diffraction peaks unique to LiNH_2 (e.g., $2\theta \approx 17^\circ$ and 20° in Fig. 1) and the decrease in intensity of diffraction peaks from LiH . Finally, the Li_3N peaks we observe during de-hydriding are due to unreacted material remaining from the initial hydriding step.

Isothermal volumetric measurements of the hydriding and de-hydriding of Li_3N are shown in Fig. 2. The initial hydriding measurements of Li_3N typically show non-equilibrium effects and do not represent true equilibrium pressure-composition isotherms. Although the data acquisition program waits until the pressures and temperatures do not change significantly (i.e., until the H_2 gas is in equilibrium with the apparatus), that time interval is not long enough to ensure that the sample is in thermodynamic equilibrium

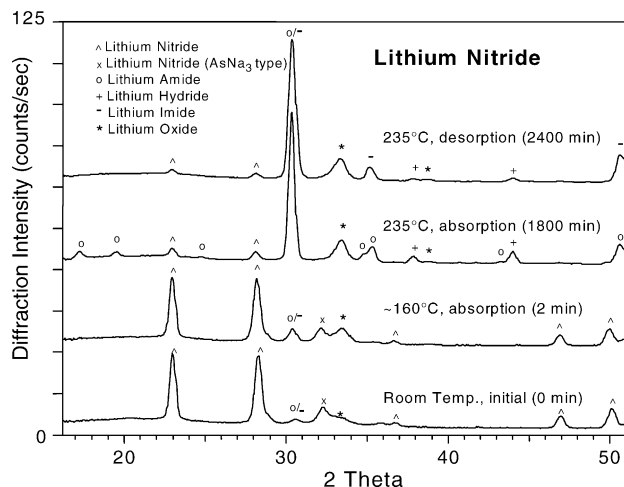


Fig. 1. Representative in situ X-ray diffraction data for hydriding and de-hydriding Li_3N .

with the H_2 gas at 1500 kPa and 159°C . This is evident by the hydriding curve decreasing as the H_2 content increases (■ in Fig. 2) and reaching a minimum near 3 wt.%. We also observed the same non-equilibrium behavior in the initial hydriding curve of another sample measured at 164°C (not shown). Hydriding is thus quite slow at these temperatures. We observe no de-hydriding into vacuum, however, implying that the kinetics of H_2 release is extremely slow. A second hydriding and de-hydriding cycle at 159°C (Δ in Fig. 2) shows only small additional H_2 uptake and, again, no H_2 release. The total H_2 absorbed by Li_3N in this experiment is $(\text{mass H}_2)/(\text{mass Li}_3\text{N} + \text{mass H}_2) = 7.2$ wt.%. The total expected H_2 content from Eq. (1) is 5.47 wt.% and Eqs. (1) and (2) is 10.37 wt.% relative to the initial starting weight of the Li_3N sample. Thus, the hydriding reaction is not complete in the time allowed for the 159°C experiment (first run: 83 h, second run: 25 h). The sample was then de-hydrided at 240°C in dynamically pumped vacuum for ~ 12 h. A third hydriding cycle (\circ in Fig. 2) indicates that the amount of cyclable H_2 at 240°C is 5.2 wt.% relative to the initial Li_3N sample weight. This result is smaller than the total H_2 absorbed previously,

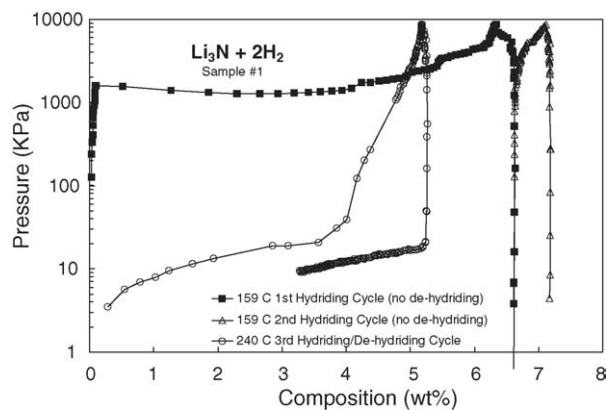


Fig. 2. Isothermal volumetric measurement of the hydriding and de-hydriding Li_3N plotted vs. log pressure.

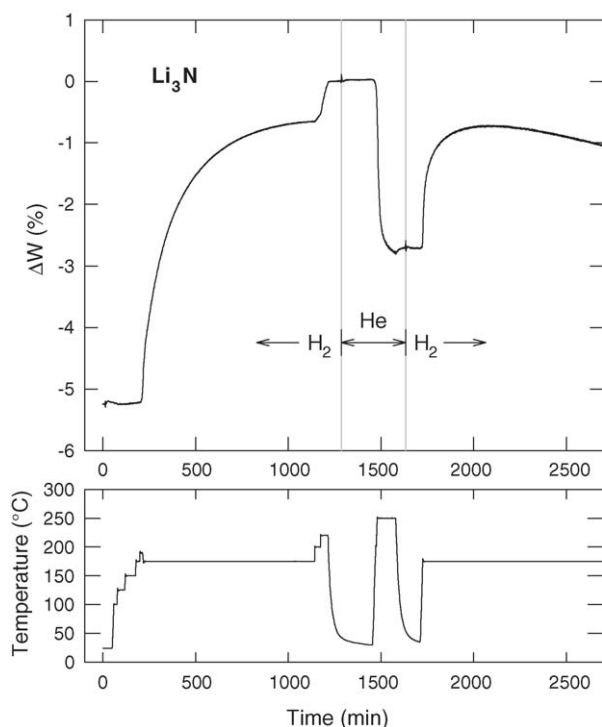


Fig. 3. Weight change with time during H_2 cycling of Li_3N . The lower panel shows the temperature profile.

but it is consistent with the reversible reaction described by Eq. (2), namely: $(\text{mass of cyclable } H_2)/(\text{mass } Li_3N + \text{mass of cyclable } H_2) = 5.47 \text{ wt.}\%$. The plateau pressures obtained are quite small, $\sim 10 \text{ kPa}$ (see Fig. 2), so that de-hydrating is inhibited when only a small H_2 pressure develops in the sample chamber.

Gravimetric measurements of the hydriding and de-hydrating cycle of Li_3N are shown in Fig. 3, where ΔW is measured relative to the weight in the hydrided state. Starting from Li_3N in flowing H_2 at 136 kPa, H_2 absorption begins at $\sim 175^\circ\text{C}$ where the sorption kinetics are slow: $\sim 4.5 \text{ wt.}\%$ H_2 uptake in $\sim 1000 \text{ min}$. Further increasing the temperature boosts H_2 absorption to 5.2 wt.% at 220°C . After cooling back to low temperature and changing the gas flow to He at 136 kPa, H_2 is released upon heating to 250°C . The total quantity of H_2 released is 2.8 wt.%, or just over half of the original quantity absorbed. This is again consistent with the conclusion that the initial weight gain corresponds to the

two-step reaction described by Eqs. (1) and (2), whereas de-hydrating under these conditions corresponds to only Eq. (2).

The sample re-hydrides after returning the gas flow to H_2 at 136 kPa and then heating to 175°C . The weight reaches a peak after regaining about 2 wt.%, then decreases again and ultimately reaches an approximately constant rate of weight loss. This weight loss is due to production of ammonia (NH_3) (see, e.g., Ref. [3]). The peak near 2100 min (and presumably also near 1100 min) occurs when the rate of H_2 uptake equals the rate of NH_3 production. We emphasize that the rate of NH_3 production is very slow; the loss of $\sim 1 \text{ wt.}\%$ over $\sim 1000 \text{ min}$ in H_2 represents the conversion of about 2% of the sample. We find that decomposition stops when the sample is below $\sim 150^\circ\text{C}$, where decomposition by NH_3 release is kinetically excluded.

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References

- [1] F.E. Pinkerton, B.G. Wicke, *Ind. Phys.* 10 (2004) 20.
- [2] P. Chen, Z. Xiong, J. Luo, J. Lin, T.K. Tan, *Nature* 420 (2002) 302.
- [3] Y. Hu, E. Ruckenstein, *J. Phys. Chem. A* 107 (2003) 9738.
- [4] Y. Hu, E. Ruckenstein, *Ind. Eng. Chem. Res.* 42 (2003) 5135.
- [5] T. Ichikawa, S. Isobe, N. Hanada, H. Fujii, *J. Alloys Compd.* 365 (2004) 271.
- [6] T. Ichikawa, N. Hanada, S. Isobe, H. Leng, H. Fujii, *J. Phys. Chem. B* 108 (2004) 7887.
- [7] Y. Nakamori, S. Orimo, *J. Alloys Compd.* 370 (2004) 271.
- [8] M.P. Balogh, G.G. Tibbetts, F.E. Pinkerton, G.P. Meisner, C.H. Olk, *J. Alloys Compd.* 350 (2003) 136.
- [9] G.G. Tibbetts, G.P. Meisner, C.H. Olk, *Carbon* 39 (2001) 2291.
- [10] G.P. Meisner, G.G. Tibbetts, F.E. Pinkerton, C.H. Olk, M.P. Balogh, *J. Alloys Compd.* 337 (2002) 254.
- [11] F.E. Pinkerton, B.G. Wicke, C.H. Olk, G.G. Tibbetts, G.P. Meisner, M.S. Meyer, J.F. Herbst, *J. Phys. Chem. B* 104 (2000) 9460.
- [12] F.E. Pinkerton, M.S. Meyer, G.G. Tibbetts, R. Chahine, *Proceedings of the 11th Canadian Hydrogen Conference, Victoria, BC, Canada, (2001) pp. 633–642.*